# Continua of Central Configurations with a Negative Mass in the $n$-Body Problem 

John B. Little

Department of Mathematics and Computer Science
College of the Holy Cross, Worcester MA
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- Julian Hachmeister (undergraduate, UH Hilo)
- Jasmine McGhee (undergraduate, Loyola Marymount)
- Roberto Pelayo (UH Hilo)
- Spencer Sasarita (undergraduate, U Arizona)

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- A c.c. (with center of mass at $\overline{\mathbf{q}}$ ) is a configuration such that the acceleration vector of each body is proportional to the displacement vector from center of mass, all with the same proportionality constant. That is,
- (Setting $G=1$ and writing $\mathbf{q}_{i}$ for location of $i$ th body):

$$
\mathbf{A}_{i}=\sum_{j \neq i} \frac{m_{j}\left(\mathbf{q}_{j}-\mathbf{q}_{i}\right)}{r_{i j}^{3}}=\omega^{2}\left(\overline{\mathbf{q}}-\mathbf{q}_{i}\right)
$$

## The big problem

- A major question here is: Given $n$ masses $m_{1}, \ldots, m_{n}$, at how many different locations can these be placed to get central configurations? (Usually in $\mathbb{R}^{2}$ or $\mathbb{R}^{3}$, but makes sense mathematically in higher dimensions too.)


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- Can translate, rotate, and scale c.c.'s and the results are again c.c.'s
- Convention: Two c.c.'s are equivalent if one can be taken into the other by a composition of a rigid motion (translation, rotation) and a scaling in $\mathbb{R}^{k}$
- More precise form of question: Is the set of equivalence classes of (planar, or ... ) central configurations finite? On Smale's 21st century problem list.


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- Only fairly limited special cases known in general
- Question is subtle algebraically. For instance, by work of Gareth Roberts (Physica D 127 (1999), 141-145), there collections of $n=5$ masses, one negative, for which there is a curve of equivalence classes of c.c.'s (a "continuum")


## Geometry of Roberts' "rhombus +1"

- Choose coordinates so

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\mathbf{q}_{0}=(0,0), \mathbf{q}_{1}=(\cos (t), 0)=-\mathbf{q}_{2}, \mathbf{q}_{3}=(0, \sin (t))=-\mathbf{q}_{4} .
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- The center of mass of the configuration, $\overline{\mathbf{q}}$, is located at the origin.


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Figure: Rhombus with Roberts' parametrization

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\begin{aligned}
\mathbf{A}_{3, y} & =-\frac{-\sin (\theta)}{4 \sin ^{3}(\theta)}-\sin (\theta)-\sin (\theta)+\frac{2 \sin (\theta)}{8 \sin ^{3}(\theta)} \\
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- Therefore

$$
\mathbf{A}_{3}=(0,-2 \sin (\theta))=2(0,-\sin (\theta))=2\left(\overline{\mathbf{q}}-\mathbf{q}_{3}\right) .
$$

## How it works, cont.

- The accelerations for each of the other bodies are similar: $\mathbf{A}_{i}=2\left(\overline{\mathbf{q}}-\mathbf{q}_{i}\right)$ for each $i=1, \ldots, 4$.


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- $\overline{\mathbf{q}}$ is fixed at the origin and the distances from the 0th body are changing but the distances between consecutive vertices of the rhombus are not
- Therefore, we have found a continuum of inequivalent c.c.'s one for each $\theta$ in the interval $0<\theta<\frac{\pi}{2}$.


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- They came up with a beautiful construction and a whole infinite family of additional examples, but only in $\mathbb{R}^{2 k}$ for $k \geq 2$.
- The found their examples by looking at Roberts' construction in a different way (but can also make them look similar, and that's what we'll do in this talk)


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- The acceleration on body 3 due to those other two masses is:

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\mathbf{A}_{3} & =\frac{m_{0}}{r_{03}^{3}}\left(\mathbf{q}_{0}-\mathbf{q}_{3}\right)+\frac{m_{4}}{r_{34}^{3}}\left(\mathbf{q}_{4}-\mathbf{q}_{3}\right) \\
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- $\mathbf{A}_{0}, \mathbf{A}_{4}$ in this sub-configuration also zero.
- Similarly for other sub-configuration $\left\{\mathbf{q}_{1}, \mathbf{q}_{0}, \mathbf{q}_{2}\right\}$.


## Neutral configurations

## Definition 1

We will say a configuration of $\ell>1$ bodies is neutral if the gravitational acceleration on each body is zero.

Easy to see that neutral configurations are only possible if at least one mass is negative.

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- Place a new mass $m_{0}$ at the center of mass $\Rightarrow$ still central
- Can take $\mathbf{q}_{0}=(0,0)$ and

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\mathbf{q}_{j}=\left(\cos \left(\frac{2 \pi j}{n}\right), \sin \left(\frac{2 \pi j}{n}\right)\right)
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for $j=1, \ldots, n$.

- Because of the symmetry, $\mathbf{A}_{n}$ has $y$-component $=0$ and

$$
\mathbf{A}_{n, x}=-m_{0}+\sum_{j=1}^{n-1} \frac{\cos \left(\frac{2 \pi j}{n}\right)-1}{\left(2-2 \cos \left(\frac{2 \pi j}{n}\right)\right)^{\frac{3}{2}}} .
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m_{0}=-\frac{\sqrt{-\sqrt{5}+5} \sqrt{2}+\sqrt{\sqrt{5}+5} \sqrt{2}}{2 \sqrt{5}}
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- With this $m_{0}$, the ( $n$-gon) +1 configuration becomes a neutral configuration because of the rotational symmetry.


## A general result

## Definition 2


#### Abstract

A $k$-dimensional regular polytope configuration is a configuration $\mathcal{C}$ of equal masses located at the vertices of a regular polytope $\mathcal{P}$ in $\mathbb{R}^{k}$ such that $\mathcal{P}$ that is not contained in any hyperplane.


## Theorem 3

There exists a negative mass $m_{0}$ that, when placed at the center of mass of a regular polytope configuration $\mathcal{C}$, creates a neutral configuration, $\mathcal{C}_{0}$.

In fact $m_{0}$ is $-\omega^{2}$ from the configuration $\mathcal{P}$.

## A generalization

## Theorem 4

Let $\mathcal{C}$ be any union of congruent regular polytope configurations in orthogonal subspaces in $\mathbb{R}^{k}$, all with center of mass at the origin. There exists a negative mass which, placed at the origin, makes the configuration $\mathcal{C}_{0}=\mathcal{C} \cup\{\mathbf{0}\}$ neutral.

- Thanks to my colleague at Holy Cross, Andy Hwang, for suggesting this idea.


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- Thanks to my colleague at Holy Cross, Andy Hwang, for suggesting this idea.
- Also, some experiments I have done indicate that the hypothesis of congruence is not necessary, if the masses in each regular polytope configuration can be different.


## Aside on regular polytopes

There is a complete classification of the regular polytopes in $\mathbb{R}^{k}$ up to similarity (see the classic book by Coxeter):
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(3) There are 6 regular polytopes in $\mathbb{R}^{4}$
(1) There are 3 regular polytopes in $\mathbb{R}^{k}, k \geq 5$ (simplex, hypercube, cross-polytope)

## A general construction

## Definition 5

Given a configuration $\mathcal{C}$ in $\mathbb{R}^{k}$, the doubling of $\mathcal{C}$ is the parametrized family of configurations for $\theta \in\left(0, \frac{\pi}{2}\right)$ in $\mathbb{R}$ defined by:

$$
\begin{aligned}
\mathcal{D}_{\theta}(\mathcal{C})= & \left\{(\cos (\theta) \mathbf{q}, \mathbf{0}) \in \mathbb{R}^{2 k}: \mathbf{q} \in \mathcal{C}\right\} \\
& \cup\left\{(\mathbf{0}, \mathbf{0}) \in \mathbb{R}^{2 k}\right\} \\
& \cup\left\{(\mathbf{0}, \sin (\theta) \mathbf{q}) \in \mathbb{R}^{2 k}: \mathbf{q} \in \mathcal{C}\right\}
\end{aligned}
$$

## How we apply this

Consider this situation:
(1) $\mathcal{C}$ is a $k$-dimensional regular polytope configuration, with $n=$ number of vertices of the polytope, vertices $\mathbf{q}_{i}$ with $\left\|\mathbf{q}_{i}\right\|=1$, all $i$ (or one of the more general configurations from Theorem 4, with $n=$ total number of vertices), all masses $=1$

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(2) $\mathcal{C}_{0}=\mathcal{C} \cup\{\mathbf{0}\}$ is an associated neutral configuration, and
(3) The $\mathcal{D}_{\theta}(\mathcal{C})$ are $(2 n+1)$-body configurations, with all masses $=1$ except for the central negative mass

## The theorem

## Theorem 6

Let $\mathcal{C}$ be a $k$-dimensional regular polytope configuration, or one of the more general "product regular polytope configurations" given in Theorem 4 in $\mathbb{R}^{k}$. Let $n$ be the number of bodies in $\mathcal{C}$. Let $\mathcal{C}_{0}$ be the associated neutral configuration. Then for each $\theta \in\left(0, \frac{\pi}{2}\right), \mathcal{D}_{\theta}(\mathcal{C})$ is a central configuration with $\omega^{2}=n$.

## Corollary 7

The family $\mathcal{D}_{\theta}(\mathcal{C})$ is a continuum of inequivalent central configurations in $\mathbb{R}^{2 k}$, all with the same masses.

## Idea of proof

- The proof is a direct check that the c.c. conditions are satisfied for each body in the doubled configuration.
- Symmetry is used in a crucial way to simplify the calculations
- What really makes this work is that the orthogonality of the two copies of $\mathbb{R}^{k}$ implies

$$
\left\|\left(\cos (\theta) \mathbf{q}_{i}, \mathbf{0}\right)-\left(\mathbf{0}, \sin (\theta) \mathbf{q}_{j}\right)\right\|=\sqrt{\cos ^{2}(\theta)+\sin ^{2}(\theta)}=1
$$

for all $i, j$.

## Comments

- As before, it is easy to see that $\mathcal{D}_{\theta_{1}}(\mathcal{C})$ and $\mathcal{D}_{\theta_{2}}(\mathcal{C})$ are not equivalent if $0<\theta_{1}<\theta_{2}<\frac{\pi}{2}$.


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- In our paper, we write the continuum using a different parametrization for the doubling construction that fixes the first copy and makes $\omega^{2}=\frac{n}{\left(1+t^{2}\right)^{3 / 2}}$. Equivalent to what we said here, though.


## Mahalo for your attention!

