Continua of Central Configurations with a Negative Mass in the *n*-Body Problem

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AMS Special Session on Celestial Mechanics JMM, San Diego

January 9, 2013

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Work done at PURE Math 2012 REU-type program at University of Hawai'i at Hilo; joint with

- Julian Hachmeister (undergraduate, UH Hilo)
- Jasmine McGhee (undergraduate, Loyola Marymount)
- Roberto Pelayo (UH Hilo)
- Spencer Sasarita (undergraduate, U Arizona)

To appear, Celestial Mechanics and Dynamical Astronomy.

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- (Setting G = 1 and writing \mathbf{q}_i for location of *i*th body):

$$\mathbf{A}_i = \sum_{j \neq i} \frac{m_j(\mathbf{q}_j - \mathbf{q}_i)}{r_{ij}^3} = \omega^2(\overline{\mathbf{q}} - \mathbf{q}_i)$$

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The big problem

A major question here is: Given *n* masses *m*₁,..., *m_n*, at how many different locations can these be placed to get central configurations? (Usually in ℝ² or ℝ³, but makes sense mathematically in higher dimensions too.)

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- Convention: Two c.c.'s are equivalent if one can be taken into the other by a composition of a rigid motion (translation, rotation) and a scaling in R^k
- More precise form of question: Is the set of equivalence classes of (planar, or ...) central configurations *finite*? On Smale's 21st century problem list.

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- Quite strong results for n = 5 (Albouy and Kaloshin)
- Only fairly limited special cases known in general
- Question is subtle algebraically. For instance, by work of Gareth Roberts (Physica D 127 (1999), 141-145), there collections of n = 5 masses, one negative, for which there is a *curve* of equivalence classes of c.c.'s (a "continuum")

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Geometry of Roberts' "rhombus +1"

Choose coordinates so

$$\mathbf{q}_0 = (0,0), \mathbf{q}_1 = (\cos(t),0) = -\mathbf{q}_2, \mathbf{q}_3 = (0,\sin(t)) = -\mathbf{q}_4.$$

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- Let $m_i = 1$ for i = 1, ..., 4, and $m_0 = -\frac{1}{4}$.
- The center of mass of the configuration, q
 is located at the origin.

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Roberts' "rhombus +1"

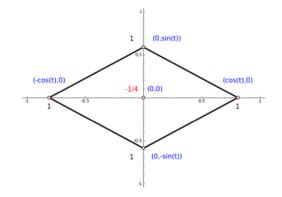


Figure: Rhombus with Roberts' parametrization

John B. Little Continua of Central Configurations

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How it works

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- The y-component is

$$\begin{aligned} \mathbf{A}_{3,y} &= -\frac{-\sin(\theta)}{4\sin^3(\theta)} - \sin(\theta) - \sin(\theta) + \frac{2\sin(\theta)}{8\sin^3(\theta)} \\ &= -2\sin(\theta). \end{aligned}$$

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Therefore

$$\mathbf{A}_3 = (\mathbf{0}, -2\sin(\theta)) = \mathbf{2}(\mathbf{0}, -\sin(\theta)) = \mathbf{2}(\overline{\mathbf{q}} - \mathbf{q}_3).$$

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How it works, cont.

• The accelerations for each of the other bodies are similar: $\mathbf{A}_i = 2(\overline{\mathbf{q}} - \mathbf{q}_i)$ for each i = 1, ..., 4.

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- From the symmetry of the configuration, the acceleration of the body at the origin cancels to **0**.
- The c.c. equations are satisfied for each θ , $0 < \theta < \frac{\pi}{2}$, with $\omega^2 = 2$.
- **q** is fixed at the origin and the distances from the 0th body are changing but the distances between consecutive vertices of the rhombus are not
- Therefore, we have found a continuum of inequivalent c.c.'s one for each θ in the interval 0 < θ < π/2.

Rethinking Roberts

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- Last summer in Hawai'i, I asked one of my REU groups to try to see whether they could find other similar examples.
- They came up with a beautiful construction and a whole infinite family of additional examples, *but* only in ℝ^{2k} for k ≥ 2.
- The found their examples by looking at Roberts' construction in a different way (but can also make them look similar, and that's what we'll do in this talk)

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An observation

 Consider the sub-configuration {q₃, q₀, q₄}, disregarding the other two bodies.

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- The acceleration on body 3 due to those other two masses is:

$$\begin{aligned} \mathbf{A}_{3} &= \frac{m_{0}}{r_{03}^{3}} \left(\mathbf{q}_{0} - \mathbf{q}_{3} \right) + \frac{m_{4}}{r_{34}^{3}} \left(\mathbf{q}_{4} - \mathbf{q}_{3} \right) \\ &= \frac{1}{4 \sin^{3}(\theta)} (0, -\sin(\theta)) + \frac{1}{8 \sin^{3}(\theta)} (0, 2\sin(\theta)) \\ &= \mathbf{0} \end{aligned}$$

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- **A**₀, **A**₄ in this sub-configuration also zero.
- Similarly for other sub-configuration $\{q_1, q_0, q_2\}$.

Neutral configurations

Definition 1

We will say a configuration of $\ell > 1$ bodies is **neutral** if the gravitational acceleration on each body is zero.

Easy to see that neutral configurations are only possible if at least one mass is negative.

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Other neutral configurations

• First ingredient is an "especially symmetric" central configuration.

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- Example regular *n*-gon configurations in the *xy*-plane with positive masses m_i = 1 are central.

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- Can take $\mathbf{q}_0 = (0,0)$ and

$$\mathbf{q}_j = \left(\cos\left(\frac{2\pi j}{n}\right), \sin\left(\frac{2\pi j}{n}\right)\right)$$

for *j* = 1, . . . , *n*.

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$$\mathbf{q}_j = \left(\cos\left(\frac{2\pi j}{n}\right), \sin\left(\frac{2\pi j}{n}\right)\right)$$

for j = 1, ..., n.

• Because of the symmetry, A_n has y-component = 0 and

$$\mathbf{A}_{n,x} = -m_0 + \sum_{j=1}^{n-1} \frac{\cos(\frac{2\pi j}{n}) - 1}{(2 - 2\cos(\frac{2\pi j}{n}))^{\frac{3}{2}}}$$

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Example, cont.

Setting this equal to zero, we can solve for m₀ to make the acceleration on the *n*th body equal to 0.

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Example, cont.

- Setting this equal to zero, we can solve for m₀ to make the acceleration on the *n*th body equal to 0.
- When n = 5, for instance, this yields

$$m_0 = -rac{\sqrt{-\sqrt{5}+5}\sqrt{2}+\sqrt{\sqrt{5}+5}\sqrt{2}}{2\sqrt{5}},$$

and there will be an analogous m_0 value for any other $n \ge 3$.

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Example, cont.

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 With this m₀, the (n-gon)+1 configuration becomes a neutral configuration because of the rotational symmetry.

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A general result

Definition 2

A *k*-dimensional regular polytope configuration is a configuration C of equal masses located at the vertices of a regular polytope P in \mathbb{R}^k such that P that is not contained in any hyperplane.

Theorem 3

There exists a negative mass m_0 that, when placed at the center of mass of a regular polytope configuration C, creates a neutral configuration, C_0 .

In fact m_0 is $-\omega^2$ from the configuration \mathcal{P} .

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A generalization

Theorem 4

Let C be any union of congruent regular polytope configurations in orthogonal subspaces in \mathbb{R}^k , all with center of mass at the origin. There exists a negative mass which, placed at the origin, makes the configuration $C_0 = C \cup \{\mathbf{0}\}$ neutral.

 Thanks to my colleague at Holy Cross, Andy Hwang, for suggesting this idea.

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- Thanks to my colleague at Holy Cross, Andy Hwang, for suggesting this idea.
- Also, some experiments I have done indicate that the hypothesis of congruence is not necessary, if the masses in each regular polytope configuration can be different.

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Aside on regular polytopes

There is a complete classification of the regular polytopes in \mathbb{R}^k up to similarity (see the classic book by Coxeter):

• The regular *n*-gons, $n \ge 3$ in \mathbb{R}^2 ,

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- 2 The 5 Platonic solids in \mathbb{R}^3 (tetrahedron, cube, octahedron, dodecahedron, icosahedron)
- **③** There are 6 regular polytopes in \mathbb{R}^4
- There are 3 regular polytopes in ℝ^k, k ≥ 5 (simplex, hypercube, cross-polytope)

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A general construction

Definition 5

Given a configuration C in \mathbb{R}^k , the **doubling of** C is the parametrized family of configurations for $\theta \in (0, \frac{\pi}{2})$ in \mathbb{R} defined by:

$$egin{aligned} \mathcal{D}_{ heta}(\mathcal{C}) &= \{(\cos(heta)\mathbf{q},\mathbf{0})\in\mathbb{R}^{2k}:\mathbf{q}\in\mathcal{C}\}\ &\cup\{(\mathbf{0},\mathbf{0})\in\mathbb{R}^{2k}\}\ &\cup\{(\mathbf{0},\sin(heta)\mathbf{q})\in\mathbb{R}^{2k}:\mathbf{q}\in\mathcal{C}\} \end{aligned}$$

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How we apply this

Consider this situation:

C is a k-dimensional regular polytope configuration, with n = number of vertices of the polytope, vertices q_i with ||q_i|| = 1, all i (or one of the more general configurations from Theorem 4, with n = total number of vertices), all masses = 1

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- 2 $\mathcal{C}_0 = \mathcal{C} \cup \{\boldsymbol{0}\}$ is an associated neutral configuration, and
- So The $D_{\theta}(C)$ are (2n + 1)-body configurations, with all masses = 1 except for the central negative mass

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The theorem

Theorem 6

Let C be a k-dimensional regular polytope configuration, or one of the more general "product regular polytope configurations" given in Theorem 4 in \mathbb{R}^k . Let n be the number of bodies in C. Let C_0 be the associated neutral configuration. Then for each $\theta \in (0, \frac{\pi}{2}), \mathcal{D}_{\theta}(C)$ is a central configuration with $\omega^2 = n$.

Corollary 7

The family $\mathcal{D}_{\theta}(\mathcal{C})$ is a continuum of inequivalent central configurations in \mathbb{R}^{2k} , all with the same masses.

Idea of proof

- The proof is a direct check that the c.c. conditions are satisfied for each body in the doubled configuration.
- Symmetry is used in a crucial way to simplify the calculations
- What really makes this work is that the orthogonality of the two copies of R^k implies

$$\|(\cos(\theta)\mathbf{q}_i,\mathbf{0})-(\mathbf{0},\sin(\theta)\mathbf{q}_j)\|=\sqrt{\cos^2(\theta)+\sin^2(\theta)}=1$$

for all *i*, *j*.

Comments

As before, it is easy to see that D_{θ1}(C) and D_{θ2}(C) are not equivalent if 0 < θ1 < θ2 < π/2.

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Comments

- As before, it is easy to see that D_{θ1}(C) and D_{θ2}(C) are not equivalent if 0 < θ1 < θ2 < π/2.
- In our paper, we write the continuum using a different parametrization for the doubling construction that fixes the first copy and makes $\omega^2 = \frac{n}{(1+t^2)^{3/2}}$. Equivalent to what we said here, though.

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Mahalo for your attention!

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